Computing wave functions in multichannel collisions with non-local potentials using the R-matrix method

Motivation

To describe quantum properties and predict the positions of particles after collisions, we use wave functions. The wave functions for two-body nuclear collisions are also important input to understand more complicated three-body collisions, specifically nuclear transfer reactions [2].

Previous solutions to this problem have been computationally time-consuming. Since wave functions appear often in nuclear physics, there is a demand for faster methods for calculating them.



The R-matrix method makes use of the assumption that nuclear forces are short-ranged. The problem is divided into two regions, internal and external, which are separated at the channel radius a. Nuclear forces are only considered in the internal region so that the wave function in the external region is calculated analytically.

Theory [1]

$$\mathcal{L}_{c}(B) = \frac{\hbar^{2}}{2\mu_{c}}\delta(r-\alpha)\left(\frac{d}{dr} - \frac{B_{c}}{r}\right)$$

First, we introduce the Bloch surface operator into the Schrödinger equation to relate the internal and external wave functions

$$[(T_c + \mathcal{L}_c + E_c - E)\delta_{cc'} + V_{cc'}]u_{c'}^{int} = \mathcal{L}_c u_c^{ext}$$
N

 $u_c^{int}(r) = \sum c_{cj} \varphi_j(r)$

Next, we expand the partial wave functions over a finite basis (see Lagrange Mesh section)

$$R_{cc'} = \frac{\hbar^2}{2\sqrt{\mu_c\mu_{c'}a}} \sum_{i,i'=1}^{N} \varphi_i(a)(\mathbf{C}^{-1})_{ci,c'i'}\varphi_{i'}(a)$$

$$C_{ci,c'i'} = \langle \varphi_i | T_c + \mathcal{L}_c + E_c - E | \varphi_{i'} \rangle \,\delta_{cc'} + \langle \varphi_i | V_{cc'} | \varphi_{i'} \rangle$$

$$u_{c(c_0)}^{int}(r) = \sum_{c'} \mathcal{L}_{c'} u_{c'(c_0)}^{ext}(a) \sum_{i,i'=1}^{N} \varphi_i(r)(C^{-1})_{ci,c'i'}\varphi_{i'}(a)$$

Lagrange Mesh

For the finite basis expansion required by the R-matrix method, we use the Lagrange functions defined here:

$$\varphi_{i}(r) = (-1)^{N+i} \left(\frac{r}{ax_{i}}\right) \sqrt{ax_{i}(1-x_{i})} \frac{P_{N}(2r/a-1)}{r-ax_{i}} \qquad \qquad \langle \varphi_{i}|\varphi_{j}\rangle = \int_{0}^{a} \varphi_{i}(r)\varphi_{j}(r)dr \approx \delta_{ij}$$

Due to the Gauss-Legendre quadrature rule, when evaluating the potentials on these functions, one only has to consider the value of the potentials at the mesh points, ax_i This makes the R-matrix method **much faster** than other techniques which use direct integration [3].





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Nuclear potentials in reality are *nonlocal*, meaning the nuclear force on a particle at one point depends on the force on that particle at all other points.

Most problems also involve multiple states, or *channels*. This introduces coupling into the equations.

 $\left[-\frac{\hbar^2}{2\mu_c}\left(\frac{d^2}{dr^2}-\frac{l_c(l_c+1)}{r^2}\right)+V_c(r)+E_c-E\right]u_{c(c_0)}(r)+\sum_{c'}\int_0^\infty V_{cc'}(r,r')u_{c'(c_0)}(r')dr'=0$

By taking the operators from the Schrödinger equation on the basis functions, we calculate the R-matrix, which is used to compute the external wave functions. We can then compute the internal wave functions at all points.

Gauss quadrature rules:

 $\langle \varphi_i | V(\mathbf{r}) | \varphi_j \rangle \approx V(\mathbf{a} \mathbf{x}_i) \delta_{ij}$

 $\langle \varphi_i | W(\mathbf{r}, \mathbf{r}') | \varphi_j \rangle \approx a \sqrt{\lambda_i \lambda_j} W(ax_i, ax_j)$





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